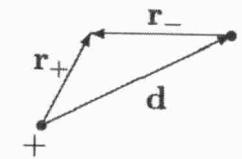


PROBLEM 2.18

From Prob. 2.12, the field inside the positive sphere is $\mathbf{E}_+ = \frac{\rho}{3\epsilon_0} \mathbf{r}_+$, where \mathbf{r}_+ is the vector from the positive center to the point in question. Likewise, the field of the negative sphere is $-\frac{\rho}{3\epsilon_0} \mathbf{r}_-$. So the *total* field is

$$\mathbf{E} = \frac{\rho}{3\epsilon_0} (\mathbf{r}_+ - \mathbf{r}_-)$$

But (see diagram) $\mathbf{r}_+ - \mathbf{r}_- = \mathbf{d}$. So $\boxed{\mathbf{E} = \frac{\rho}{3\epsilon_0} \mathbf{d}}$



Problem 2.20

$$(1) \nabla \times \mathbf{E}_1 = k \begin{vmatrix} \hat{\mathbf{x}} & \hat{\mathbf{y}} & \hat{\mathbf{z}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ xy & 2yz & 3zx \end{vmatrix} = k [\hat{\mathbf{x}}(0 - 2y) + \hat{\mathbf{y}}(0 - 3z) + \hat{\mathbf{z}}(0 - x)] \neq \mathbf{0},$$

so \mathbf{E}_1 is an *impossible* electrostatic field.

$$(2) \nabla \times \mathbf{E}_2 = k \begin{vmatrix} \hat{\mathbf{x}} & \hat{\mathbf{y}} & \hat{\mathbf{z}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ y^2 & 2xy + z^2 & 2yz \end{vmatrix} = k [\hat{\mathbf{x}}(2z - 2z) + \hat{\mathbf{y}}(0 - 0) + \hat{\mathbf{z}}(2y - 2y)] = \mathbf{0},$$

so \mathbf{E}_2 is a *possible* electrostatic field.

Let's go by the indicated path:

$$\mathbf{E} \cdot d\mathbf{l} = (y^2 dx + (2xy + z^2)dy + 2yz dz)k$$

Step I: $y = z = 0$; $dy = dz = 0$. $\mathbf{E} \cdot d\mathbf{l} = ky^2 dx = 0$.

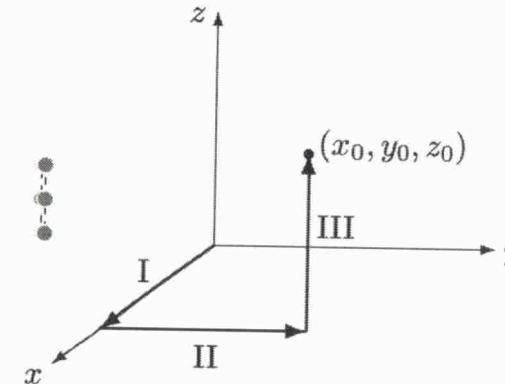
Step II: $x = x_0$, $y : 0 \rightarrow y_0$, $z = 0$. $dx = dz = 0$.

$$\mathbf{E} \cdot d\mathbf{l} = k(2xy + z^2)dy = 2kx_0 y dy.$$

$$\int_{II} \mathbf{E} \cdot d\mathbf{l} = 2kx_0 \int_0^{y_0} y dy = kx_0 y_0^2.$$

Step III: $x = x_0$, $y = y_0$, $z : 0 \rightarrow z_0$; $dx = dy = 0$.

$$\mathbf{E} \cdot d\mathbf{l} = 2ky_0 z dz = 2ky_0 z dz.$$



$$\int_{III} \mathbf{E} \cdot d\mathbf{l} = 2y_0 k \int_0^{z_0} z dz = ky_0 z_0^2.$$

$$V(x_0, y_0, z_0) = - \int_0^{\infty} \mathbf{E} \cdot d\mathbf{l} = -k(x_0 y_0^2 + y_0 z_0^2), \text{ or } \boxed{V(x, y, z) = -k(xy^2 + yz^2)}.$$

Check: $-\nabla V = k \left[\frac{\partial}{\partial x} (xy^2 + yz^2) \hat{\mathbf{x}} + \frac{\partial}{\partial y} (xy^2 + yz^2) \hat{\mathbf{y}} + \frac{\partial}{\partial z} (xy^2 + yz^2) \hat{\mathbf{z}} \right] = k [y^2 \hat{\mathbf{x}} + (2xy + z^2) \hat{\mathbf{y}} + 2yz \hat{\mathbf{z}}] = \mathbf{E}$. ✓

PROBLEM 2.34

(a) $W = \frac{1}{2} \int \rho V d\tau$. From Prob. 2.21 (or Prob. 2.28): $V = \frac{\rho}{2\epsilon_0} \left(R^2 - \frac{r^2}{3} \right) = \frac{1}{4\pi\epsilon_0} \frac{q}{2R} \left(3 - \frac{r^2}{R^2} \right)$

$$\begin{aligned} W &= \frac{1}{2} \rho \frac{1}{4\pi\epsilon_0} \frac{q}{2R} \int_0^R \left(3 - \frac{r^2}{R^2} \right) 4\pi r^2 dr = \frac{q\rho}{4\epsilon_0 R} \left[3 \frac{r^3}{3} - \frac{1}{R^2} \frac{r^5}{5} \right] \Big|_0^R = \frac{q\rho}{4\epsilon_0 R} \left(R^3 - \frac{R^3}{5} \right) \\ &= \frac{q\rho}{5\epsilon_0} R^2 = \frac{qR^2}{5\epsilon_0} \frac{q}{\frac{4}{3}\pi R^3} = \boxed{\frac{1}{4\pi\epsilon_0} \left(\frac{3}{5} \frac{q^2}{R} \right)}. \end{aligned}$$

(b) $W = \frac{\epsilon_0}{2} \int E^2 d\tau$. Outside ($r > R$) $\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}$; Inside ($r < R$) $\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} r \hat{\mathbf{r}}$.

$$\begin{aligned} \therefore W &= \frac{\epsilon_0}{2} \frac{1}{(4\pi\epsilon_0)^2} q^2 \left\{ \int_R^\infty \frac{1}{r^4} (r^2 4\pi dr) + \int_0^R \left(\frac{r}{R^3} \right)^2 (4\pi r^2 dr) \right\} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \left\{ \left(-\frac{1}{r} \right) \Big|_R^\infty + \frac{1}{R^6} \left(\frac{r^5}{5} \right) \Big|_0^R \right\} = \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \left(\frac{1}{R} + \frac{1}{5R} \right) = \frac{1}{4\pi\epsilon_0} \frac{3}{5} \frac{q^2}{R}. \checkmark \end{aligned}$$

(c) $W = \frac{\epsilon_0}{2} \{ \oint_S V \mathbf{E} \cdot d\mathbf{a} + \int_V E^2 d\tau \}$, where V is large enough to enclose all the charge, but otherwise arbitrary. Let's use a sphere of radius $a > R$. Here $V = \frac{1}{4\pi\epsilon_0} \frac{q}{r}$.

$$\begin{aligned} W &= \frac{\epsilon_0}{2} \left\{ \int_{r=a} \left(\frac{1}{4\pi\epsilon_0} \frac{q}{r} \right) \left(\frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \right) r^2 \sin\theta d\theta d\phi + \int_0^R E^2 d\tau + \int_R^a \left(\frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \right)^2 (4\pi r^2 dr) \right\} \\ &= \frac{\epsilon_0}{2} \left\{ \frac{q^2}{(4\pi\epsilon_0)^2} \frac{1}{a} 4\pi + \frac{q^2}{(4\pi\epsilon_0)^2} \frac{4\pi}{5R} + \frac{1}{(4\pi\epsilon_0)^2} 4\pi q^2 \left(-\frac{1}{r} \right) \Big|_R^a \right\} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \left\{ \frac{1}{a} + \frac{1}{5R} - \frac{1}{a} + \frac{1}{R} \right\} = \frac{1}{4\pi\epsilon_0} \frac{3}{5} \frac{q^2}{R}. \checkmark \end{aligned}$$

As $a \rightarrow \infty$, the contribution from the surface integral $\left(\frac{1}{4\pi\epsilon_0} \frac{q^2}{2a} \right)$ goes to zero, while the volume integral $\left(\frac{1}{4\pi\epsilon_0} \frac{q^2}{2a} \left(\frac{6a}{5R} - 1 \right) \right)$ picks up the slack.