453 Midterm (17 points). 7 Oct pt 08, 14:50 – 16:15 p.m., rm 303

Problem 1 (5 points).

A wire loop of radius *a* and resistance *R* lies in the *XY* plane. There is a uniform magnetic field $\vec{B} = B\hat{z}$ filling the whole space. What total charge passes a given point in the loop when it is rotated by 90° around the x axis?

Solution:

The induced current is

$$I(t) = \frac{1}{R}\mathcal{E}(t) = -\frac{1}{R}\frac{d\Phi}{dt}$$

so the total charge passing a given point has the form

$$Q = \int_0^\infty I(t)dt = -\frac{1}{R} \int_0^\infty \frac{d\Phi}{dt} dt = \frac{1}{R} (\Phi_{\text{initial}} - \Phi_{\text{final}}) = \frac{\pi a^2 B}{R}$$

Problem 2 (6 points).

A long cylindrical metal wire of radius a and length $l \gg a$ carries a uniform current of density J in the axial direction. The conductivity of the metal is σ .

(a) Find Poynting vector at the surface of the wire.

(b) Calculate the flux of energy through the surface of a cylinder.

Solution

From Ampere's law

$$\vec{B}(s) \stackrel{s \ge a}{=} \frac{\mu_0 I}{2\pi s} \hat{\phi} \quad \Rightarrow \quad \vec{B}(a) = \frac{\mu_0 I}{2\pi a} \hat{\phi} = \frac{\mu_0}{2} J a \hat{\phi}$$

From Ohm's law

$$\vec{J} = \sigma \vec{E} \quad \Rightarrow \quad \vec{E} = \frac{1}{\sigma} J \hat{e}_3$$

 \Rightarrow the Poynting vector is

$$\vec{S}(a) = \frac{1}{\mu_0} \vec{E}(a) \times \vec{B}(a) = \frac{J^2 a}{2\sigma} \hat{e}_3 \times \hat{\phi} = -\frac{J^2 a}{2\sigma} \hat{e}_7$$

The sign reflects the fact that the energy is pumped *into* the wire (to be dissipated like the ohmic heat). Flux of the energy is

$$P = \int dA |\vec{S}| = 2\pi a l |\vec{S}| = \frac{\pi a^2 J^2}{\sigma} l$$

Check: ohmic heat

$$P = I^2 R = (J\pi a^2)^2 \frac{l}{\sigma \pi a^2} = \frac{\pi a^2 J^2}{\sigma} l$$

Problem 3 (5 points).

A linearly polarized electromagnetic plane wave is normally incident on an infinitely large plane made from a perfect conductor. Find the charge and current densities induced on the conducting plane.

Extra credit - 3 points. Do the same for circularly polarized electromagnetic plane

Solution

Let us choose the z axis in the direction normal to the plane and the wave coming from above the plane with polarization in x direction. The incident wave has the form $(\hat{e}_3 \times \hat{e}_+ = -i\hat{e}_+)$

$$\vec{E}_{I} = \hat{e}_{1}E_{0}e^{-i\omega t+ikz}, \quad \vec{B}_{I} = \frac{\hat{n}}{c} \times \vec{E} = -\frac{\hat{e}_{3}}{c} \times \vec{E} = -\hat{e}_{2}\frac{E_{0}}{c}e^{-i\omega t+ikz}$$

The boundary condition for the electric field is $E_{\parallel}^{\text{above}} = E_{\parallel}^{\text{below}} = 0$ so $\vec{E}_R = i\hat{e}_+E_0$ and the reflected wave is

$$\vec{E}_R = -\hat{e}_1 E_0 e^{-i\omega t - ikz}, \qquad \vec{B}_R = \frac{\hat{e}_3}{c} \times \vec{E}_R \ e^{-i\omega t - ikz} = -\hat{e}_2 \frac{E_0}{c} e^{-i\omega t - ikz}$$

The sum of the reflected and incident waves takes the form:

$$\vec{E} = 2\hat{e}_1 E_0 \sin kz \ e^{-i\omega t}, \quad \vec{B} = -2\hat{e}_2 \frac{E_0}{c} \cos kz \ e^{-i\omega t}$$

Now, σ and K can be found from the remaining boundary conditions

$$\epsilon_1 \vec{E}_{\perp}^{\text{above}} - \epsilon_2 \vec{E}_{\perp}^{\text{below}} = \sigma, \quad \frac{1}{\mu_1} \vec{B}_{\perp}^{\text{above}} - \frac{1}{\mu_2} \vec{B}_{\perp}^{\text{below}} = \vec{K} \times \hat{e}_3$$

We get $\sigma = 0$ and

$$\mu_0 \vec{K} \times \hat{e}_3 = -2\hat{e}_2 \frac{E_0}{c} \ e^{-i\omega t} \ \Rightarrow \ \hat{e}_3 \times (\vec{K} \times \hat{e}_3) = \ \vec{K} = -2\hat{e}_3 \times \hat{e}_2 \frac{E_0}{\mu_0 c} \ e^{-i\omega t} = 2\hat{e}_1 \frac{E_0}{\mu_0 c} \ e^{-i\omega t}$$

. Choosing E_0 to be real and taking real part we get

$$\vec{K} = \frac{2}{\mu_0 c} \hat{e}_1 E_0 \, \cos \omega t$$

Extra problem

Let us now consider plane wave linearly polarized in y direction. The incident and reflected waves are

$$\vec{E}_{I} = 2\hat{e}_{2}E_{0} e^{ikz-i\omega t}, \quad \vec{B}_{I} = 2\hat{e}_{1}\frac{E_{0}}{c} e^{ikz-i\omega t}$$
$$\vec{E}_{R} = -2\hat{e}_{2}E_{0} e^{-ikz-i\omega t}, \quad \vec{B}_{R} = \frac{\hat{e}_{3}}{c} \times \vec{E}_{R} e^{-i\omega t-ikz} = 2\hat{e}_{1}\frac{E_{0}}{c} e^{-ikz-i\omega t}$$

and the sum of the incident and reflected waves takes the form:

$$\vec{E} = 2\hat{e}_2 E_0 \sin kz \ e^{-i\omega t}, \quad \vec{B} = -2\hat{e}_1 \frac{E_0}{c} \cos kz \ e^{-i\omega t}$$

Same calculation yields $\sigma = 0$ and

$$\mu_0 \vec{K} \times \hat{e}_3 = 2\hat{e}_1 \frac{E_0}{c} \ e^{-i\omega t} \ \Rightarrow \ \hat{e}_3 \times (\vec{K} \times \hat{e}_3) = \ \vec{K} = 2\hat{e}_3 \times \hat{e}_1 \frac{E_0}{\mu_0 c} \ e^{-i\omega t} = 2\hat{e}_2 \frac{E_0}{\mu_0 c} \ e^{-i\omega t}$$

. The circularly polarized wave is a sum of x-polarized wave and y-polarized wave with phase shift $\pm \frac{\pi}{2}$ so

$$\vec{K} = \frac{2}{\mu_0 c} \hat{e}_1 E_0 \cos \omega t + \frac{2}{\mu_0 c} \hat{e}_2 E_0 \cos \left(\omega t \pm \frac{\pi}{2}\right)$$

In components

$$K_x = \frac{2}{\mu_0 c} E_0 \cos \omega t, \quad K_y = \mp \frac{2}{\mu_0 c} E_0 \sin \omega t$$