Solution

The projection operators \hat{P}_0 and \hat{P}_1 can be rewritten as

$$\hat{P}_1 \ = \ \frac{3}{4} - \vec{\hat{s}}_1 \cdot \vec{\hat{s}}_2 \ = \ \frac{1}{2} \hat{s}^2, \quad \hat{P}_0 \ = \ 1 - \frac{1}{2} \hat{s}^2 \quad \Rightarrow \quad V \ = \ [v_c(r) - \frac{3}{4} v_\sigma(r)] \hat{P}_0 + [v_c(r) + \frac{1}{4} v_\sigma(r)] \hat{P}_1$$

Since $\hat{s}^2|1,m\rangle = 2|1,m\rangle$ and $\hat{s}^2|0,m=0\rangle = 0$ it is clear that

$$\hat{P}_1|1,m\rangle = |1,m\rangle, \qquad \hat{P}_1|0,m=0\rangle = 0$$

and

$$\hat{P}_0|0, m = 0\rangle = |0, m = 0\rangle, \qquad \hat{P}_0|1, m\rangle = 0.$$

The Hamiltonian of "relative particle" is

$$\hat{H} \; = \; \frac{\hat{p}^2}{2\mu} - E + (v_c - \frac{3}{4}v_\sigma)\hat{P}_0 + (v_c + \frac{1}{4}v_\sigma)\hat{P}_1 \; = \; \frac{p^2}{2\mu} - E + [v_c(r) - \frac{3}{4}v_\sigma(r)](1 - \frac{\hat{s}^2}{2}) + [v_c(r) + \frac{1}{4}v_\sigma(r)]\frac{\hat{s}^2}{2}$$

Schrödinger equation:

$$\begin{split} &(\hat{H}-E)|\psi_{E}\rangle \ = \ \lambda_{00}\Big[\frac{\hat{p}^{2}}{2\mu}-E+[v_{c}(r)-\frac{3}{4}v_{\sigma}(r)](1-\frac{\hat{s}^{2}}{2})+[v_{c}(r)+\frac{1}{4}v_{\sigma}(r)]\frac{\hat{s}^{2}}{2}\Big]|\psi_{E}^{(0)}\rangle|s=0,m=0\rangle \\ &+\ \sum_{m=-1}^{m=1}\lambda_{1m}\Big[\frac{\hat{p}^{2}}{2\mu}-E+[v_{c}(r)-\frac{3}{4}v_{\sigma}(r)](1-\frac{\hat{s}^{2}}{2})+[v_{c}(r)+\frac{1}{4}v_{\sigma}(r)]\frac{\hat{s}^{2}}{2}\Big]|\psi_{E}^{(1)}\rangle|s=1,m\rangle \ = \ 0 \end{split}$$

Since $\hat{P}_0|s=1,m\rangle = \hat{P}_1|s=0,m=0\rangle = 0$ we get

$$(\hat{H} - E)|\psi_{E}\rangle = \lambda_{00} \left[\frac{\hat{p}^{2}}{2\mu} - E + [v_{c}(r) - \frac{3}{4}v_{\sigma}(r)](1 - \frac{\hat{s}^{2}}{2}) \right] |\psi_{E}^{(0)}\rangle |s = 0, m = 0 \rangle$$

$$+ \sum_{m=-1}^{m=1} \lambda_{1m} \left[\frac{\hat{p}^{2}}{2\mu} + [v_{c}(r) + \frac{1}{4}v_{\sigma}(r)] \frac{\hat{s}^{2}}{2} \right] |\psi_{E}^{(1)}\rangle |s = 1, m \rangle$$

$$= \lambda_{00} \left[\frac{\hat{p}^{2}}{2\mu} - E + [v_{c}(r) - \frac{3}{4}v_{\sigma}(r)] \right] |\psi_{E}^{(0)}\rangle |s = 0, m = 0 \rangle$$

$$+ \sum_{m=1}^{m=1} \lambda_{1m} \left[\frac{\hat{p}^{2}}{2\mu} + [v_{c}(r) + \frac{1}{4}v_{\sigma}(r)] \right] |\psi_{E}^{(1)}\rangle |s = 1, m \rangle = 0$$

Multiplying from the left by the bra states $\langle s=1,m|$ and $\langle s=0,m=0|$ we see that the equations for $\psi_E^{(0)}(r)$ and $\psi_E^{(1)}(r)$ are

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + \left[v_c(r) - \frac{3}{4} v_{\sigma}(r) \right] \right] \psi_E^{(0)}(r) = E \psi_E^{(0)}(r)$$

and

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + \left[v_c(r) + \frac{1}{4} v_{\sigma}(r) \right] \right] \psi_E^{(1)}(r) = E \psi_E^{(1)}(r)$$